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On the Frequency-domain Definition of the Monogenic Signal

by

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Abstract - This paper presents the frequency-domain definition of the monogenic signal. It is derived basing on the notion of the Quaternionic Fourier Transform.

Keywords: Quaternionic Fourier Transform, quaternionic and monogenic signals

I. INTRODUCTION

In a very interesting recent paper [1] the authors defined a new notion of the 2-D monogenic signal. It is derived using Riesz transforms called also "isotropic Hilbert transforms" [2]. Its signal-domain definition has the form

$$\psi_M(x_1, x_2) = u(x_1, x_2) + iv_{r1}(x_1, x_2) + jv_{r2}(x_1, x_2)$$
(1)

where *u* is a 2-D real signal, v_{r1} and v_{r2} are the corresponding Riesz transforms. They have the form of convolutions of *u* with so called Riesz kernels r_1 and r_2 , i.e.,

$$v_{r1}(x_1, x_2) = u(x_1, x_2) * * r_1(x_1, x_2), \quad v_{r2}(x_1, x_2) = u(x_1, x_2) * * r_2(x_1, x_2).$$
(2)

Riesz kernels and their 2-D Fourier transforms are given by the relations

$$r_{1}(x_{1},x_{2}) = \frac{x_{1}}{2\pi \left[\sqrt{x_{1}^{2} + x_{2}^{2}}\right]^{3}} \stackrel{2-D FT}{\iff} \frac{-jf_{1}}{\sqrt{f_{1}^{2} + f_{2}^{2}}}, \quad r_{2}(x_{1},x_{2}) = \frac{x_{2}}{2\pi \left[\sqrt{x_{1}^{2} + x_{2}^{2}}\right]^{3}} \stackrel{2-D FT}{\iff} \frac{-jf_{2}}{\sqrt{f_{1}^{2} + f_{2}^{2}}}.$$
(3)

In the chapter II we derive the frequency-domain definition of the monogenic signal using the inverse Quaternionic Fourier Transform (QFT) of a specific spectrum. The Appendix contains the definition and basic properties of the QFT.

II. THE FREQUENCY-DOMAIN DEFINITION OF THE MONOGENIC SIGNAL

A real signal *u* can be written in the form of a union of four terms: $u(x_1, x_2) = u_{ee} + u_{oe} + u_{eo} + u_{oo}$, where the subscripts "*e*" and "*o*" denote evenness or oddness of a function w.r.t. the corresponding variable. With this decomposition the 2-D Fourier transform has the form

$$U(f_1, f_2) = \iint (u_{ee} + u_{oe} + u_{eo} + u_{oo}) e^{-j2\pi(f_1x_1 + f_2x_2)} dx_1 dx_2 = U_{ee} - U_{oo} - j(U_{oe} + U_{eo})$$
(4)

where

$$U_{ee}(f_1, f_2) = \iint u_{ee}(x_1, x_2) \cos(2\pi f_1 x_1) \cos(2\pi f_2 x_2) dx_1 dx_2 = U_{ee}(-f_1, f_2),$$
(5)

$$U_{oo}(f_1, f_2) = \iint u_{oo}(x_1, x_2) \sin(2\pi f_1 x_1) \sin(2\pi f_2 x_2) dx_1 dx_2 = -U_{oo}(-f_1, f_2),$$
(6)

$$U_{oe}(f_1, f_2) = \iint u_{oe}(x_1, x_2) \sin(2\pi f_1 x_1) \cos(2\pi f_2 x_2) dx_1 dx_2 = -U_{oe}(-f_1, f_2),$$
(7)

$$U_{eo}(f_1, f_2) = \iint u_{eo}(x_1, x_2) \cos(2\pi f_1 x_1) \sin(2\pi f_2 x_2) dx_1 dx_2 = U_{eo}(-f_1, f_2).$$
(8)

The QFT of u may be calculated using the form [3]

$$U_{q}(f_{1},f_{2}) = QFT(u) = \frac{1-k}{2}U(f_{1},f_{2}) + \frac{1+k}{2}U(-f_{1},f_{2})$$
(9)

Notice that if $U(f_1, f_2) = U(-f_1, f_2)$, then $U_q = U$. The insertion of $U(f_1, f_2)$ and $U(-f_1, f_2)$ yields

$$U_{q}(f_{1},f_{2}) = U_{ee} - jU_{eo} + k(U_{oo} - jU_{oe}) = U_{ee} - iU_{oe} - jU_{eo} + kU_{oo}.$$
 (10)

The quaternionic spectrum of the first Riesz term is

$$QFT[v_{r1}] = V_{r1}(f_1, f_2) = (-iS_1) \cdot U_q$$
(11)

where $S_1 = \frac{f_1}{\sqrt{f_1^2 + f_2^2}}$. The insertion of (10) yields

$$QFT(v_{r1}) = (-iU_{ee} - U_{oe} + jU_{oo} + kU_{eo})S_1.$$
(12)

Therefore,

$$iQFT(v_{r1}) = (U_{ee} - iU_{oe} - jU_{eo} + kU_{oo})S_1 = U_q \cdot S_1.$$
(13)

The quaternionic spectrum of the second Riesz term is

$$QFT[v_{r_2}] = V_{r_2}(f_1, f_2) = U_q \cdot (-jS_2)$$
(14)

where $S_2 = \frac{f_2}{\sqrt{f_1^2 + f_2^2}}$. Note the inverse order in comparison to (11). The insertion of (10) yields

$$QFT(v_{r2}) = (-jU_{ee} - U_{eo} - iU_{oo} - kU_{oe})S_2.$$
(15)

Therefore,

$$jQFT(v_{r2}) = (U_{ee} - iU_{eo} - jU_{eo} + kU_{oo})S_2 = U_q \cdot S_2.$$
(16)

The addition of (10), (13) and (16) yields the following quaternionic spectrum of the monogenic signal

$$QFT[\psi_{M}] = \Gamma_{M}(f_{1}, f_{2}) = U_{q}(1 + S_{1} + S_{2})$$
$$= U_{q}\left[1 + \frac{f_{1} + f_{2}}{\sqrt{f_{1}^{2} + f_{2}^{2}}}\right].$$
(17)

where U_q is given by (10). Fig.1 shows an example of the spectrum (17) with $u(x_1, x_2) = u_{ee}(x_1, x_2) = \exp\left[-\pi\left(x_1^2 + x_2^2\right)\right]$ corresponding to $U_q(f_1, f_2) = U_{ee}(f_1, f_2) = \exp\left[-\pi\left(f_1^2 + f_2^2\right)\right]$. The total energy of the spectrum of Fig.1 equals 1 (twice the energy of the Gaussian signal which equals 0.5). The energy of the spectrum with the support in the first quadrant $(f_1 > 0, f_2 > 0)$ equals ≈ 0.69 and in the fourth quadrant $(f_1 < 0, f_2 < 0)$ only 0.0123.



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APPENDIX

A. 2-D Quaternionic Signals

The concept of the quaternion number was introduced by Hamilton in 1843 and defined by the formula

$$q = 1a + ib + jc + kd \tag{A.1}$$

where $a, b, c, d \in \Re$ and the products of the imaginary units *i*, *j*, *k* obey the rules presented in Table 1. These products are non-commutative, i.e., $ij \neq ji$.

TABLE 1.

THE PRODUCTS OF IMAGINARY UNITS

1	i	j	k
i	$i^2 = -1$	ij = k	<i>ik</i> = -j
j	<i>ji</i> = - <i>k</i>	$j^2 = -1$	jk =i
k	ki = j	kj =- i	$k^2 = -1$

The conjugate quaternion number is given by

$$q^* = 1a - ib - jc - kd \tag{A.2}$$

where the numerical coefficient 1 is usually omitted in notations. The norm (modulus) of q is defined as

$$|q| = \sqrt{qq^*} = \sqrt{a^2 + b^2 + c^2 + d^2}$$
 (A.3)

Replacing in (A.1) the numbers *a*, *b*, *c*, *d* by real functions $a(x_1, x_2)$, $b(x_1, x_2)$, $c(x_1, x_2)$ and $d(x_1, x_2)$, yields a quaternion-valued function of the form

$$q(x_1, x_2) = a(x_1, x_2) + ib(x_1, x_2) + jc(x_1, x_2) + kd(x_1, x_2),$$
(A.4)

B. The Quaternionic Fourier Transform

The Quaternionic Fourier Transform (*QFT*), that is, the quaternion-valued spectrum of the real signal $u(x_1, x_2)$, is defined by the integral [3], [4]

$$U_{q}(f_{1},f_{2}) = QFT[u(x_{1},x_{2})] = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-i2\pi f_{1}x_{1}}u(x_{1},x_{2})e^{-j2\pi f_{2}x_{2}}dx_{1}dx_{2}.$$
(A.5)

The first exponential uses the imaginary unit i and the second one, the unit j. Due to the noncommutativity of products of quaternions, the change of the order of functions in the integrand of (A.5) yields another quaternion valued function, which differs by signs of the terms. The *QFT* is invertible and its inverse is given by the integral

$$u(x_1, x_2) = QFT^{-1} \Big[U_q(f_1, f_2) \Big] = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{i2\pi f_1 x_1} U_q(f_1, f_2) e^{j2\pi f_2 x_2} df_1 df_2 .$$
(A.6)

C. The Quaternionic Hermitian Symmetry

The quaternionic spectrum defined by the *QFT* obeys the rules of the quaternionic Hermitian symmetry defined by the relations [3], [4], [5]

$$U_{q}(-f_{1},f_{2}) = \alpha_{j} \Big[U_{q}(f_{1},f_{2}) \Big],$$
(A.7)

$$U_{q}(-f_{1},-f_{2}) = \alpha_{k} \Big[U_{q}(f_{1},f_{2}) \Big],$$
(A.8)

$$U_{q}(f_{1},-f_{2}) = \alpha_{i} \Big[U_{q}(f_{1},f_{2}) \Big]$$
(A.9)

where the functions α_i , α_j and α_k are called *involutions* of U_q . Using (10) the involutions α_i , α_j and α_k are defined as follows:

$$\alpha_i(f_1, f_2) = -iU_q i = U_{ee} - iU_{oe} + jU_{eo} - kU_{oo}, \qquad (A.10)$$

$$\alpha_{j}(f_{1},f_{2}) = -jU_{q}j = U_{ee} + iU_{oe} - jU_{eo} - kU_{oo}, \qquad (A.11)$$

$$\alpha_k(f_1, f_2) = -kU_q k = U_{ee} + iU_{oe} + jU_{eo} + kU_{oo}.$$
(A.12)

These involutions can be easily derived using the rules of Table 1. The quaternionic Hermitian symmetry is illustrated in Fig.2. Any real signal can be reconstructed using a single-quadrant quaternionic spectrum. The spectral information in the complementary three quadrants is redundant.

Fig. 2. The quaternionic Hermitian symmetry